

3aSC2

# Physical mechanisms of phonation onset

## - influence of flow separation on phonation onset

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# Introduction

- Phonation sustained by a net energy transfer from the glottal flow into the vocal folds.
- **Mucosal wave, or Synchronization of two structural eigenmodes by the glottal flow**
  - Provide a favorable phase relationship between the intraglottal pressure and the vocal fold tissue velocity
  - Primary mechanism for energy transfer
- **Synchronization pattern (near field fluid-structure interaction) determines phonation threshold pressure, frequency, and vocal fold vibration**
- Acoustics-structure coupling with subglottal or supraglottal acoustics
  - Minor role, at onset of normal phonation



# Motivation/Objective

- Glottal flow phenomena
  - Flow separation
    - 1. flow separation may oscillate along the vocal fold surface, depending on the glottal channel geometry (convergent, straight, and divergent)
    - 2. may be asymmetric
    - 3. Coanda effect
  - Downstream of the flow separation point:
    - unsteady vortex shedding
    - jet instabilities (symmetric, asymmetric)
    - Turbulence
- Effects of these flow phenomena on phonation onset mechanisms/vocal fold vibration are still not well understood.

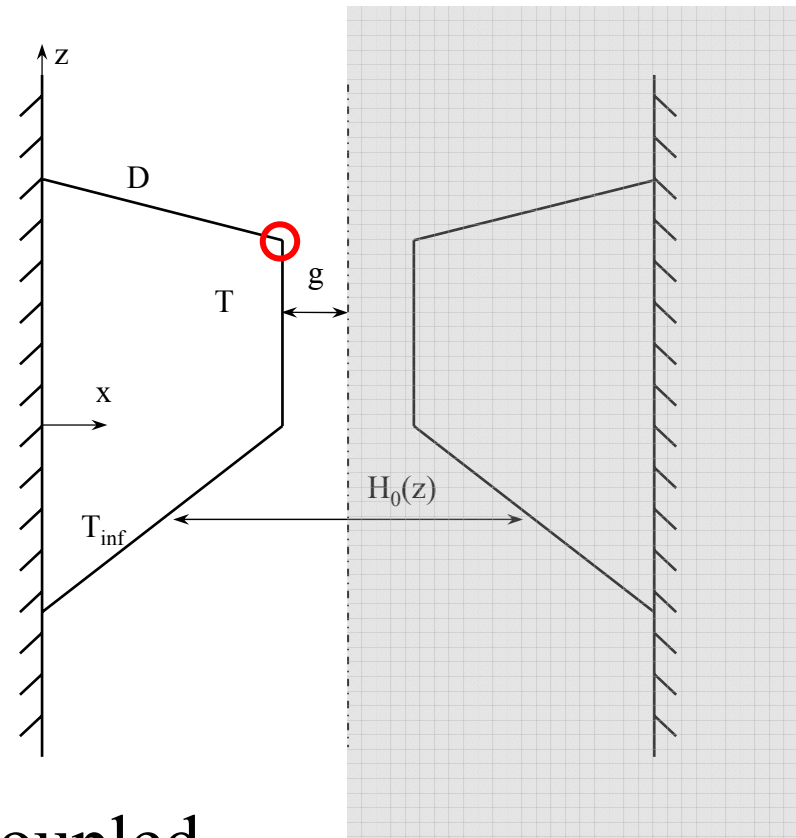


Understand the influence of flow separation on phonation onset mechanisms in terms of eigenmode-synchronization.



# Approach and Model Description

- A continuum model of the vocal fold
  - directly measurable model parameters
  - Physical description of the structure, the fluid flow, and their interaction.
  - Linear stress-strain relationship
- Flow:
  - 1D potential flow
  - Fixed separation point at superior edge
  - Neglect sub- and supra-glottal acoustic loading
- Imposed symmetry with the glottal channel centerline



Linear stability analysis of the coupled fluid-structure system.



# Non-dimensional formulation

- Length: vocal fold thickness  $T$
- Density: vocal fold density  $\rho$
- Pressure: vocal fold Young's modulus  $E$
- Velocity: wave velocity of the vocal fold structure  $\sqrt{\frac{E}{\rho_{vf}}}$
- Time:  $\frac{1}{T} \sqrt{\frac{E}{\rho_{vf}}}$
- Frequency:  $T \sqrt{\frac{\rho_{vf}}{E}}$



# Problem Formulation

- Lagrange's Equation for the vocal fold structure

$$\frac{d}{dt} \left( \frac{\partial L}{\partial \dot{q}_r} \right) - \frac{\partial L}{\partial q_r} = Q_r, r = 1, 2, \dots, N$$

- Lagrangian  $L=T-U$

$$T = \frac{1}{2} \iint_S (\dot{\xi}^2 + \dot{\eta}^2) \rho dS; \quad U = \frac{1}{2} \iint_S (\varepsilon_x \tau_x + \varepsilon_z \tau_z + \gamma_{xz} \tau_{xz}) dS$$

Displacement in x  
and z direction:

$$\xi(x, z, t) = \sum_{i=1}^I \sum_{k=0}^K A_{ik}(t) x^i z^k; \quad \eta(x, z, t) = \sum_{i=1}^I \sum_{k=0}^K B_{ik}(t) x^i z^k$$

Generalized  
Coordinates:

$$[q] = [A_{10}, A_{11}, \dots, A_{IK}, B_{10}, \dots, B_{IK}]_{N=2 \times I \times (K+1)}$$

Generalized Force:

$$Q_r = - \int_{l_{FSI}} \left( p \frac{\partial \xi}{\partial q_r} n_x + p \frac{\partial \eta}{\partial q_r} n_z \right) dl, r = 1, 2, \dots, N$$

$p$ : Flow Pressure;  $l_{FSI}$ : fluid-structure interaction interface



# Flow Model Assumptions

- One-dimensional potential flow up to the point of flow separation;
- Fixed flow separation point at the superior edge of the vocal folds.
  - This assumption is reasonable as this study concerns the linear stability of the coupled fluid-structure system.
- Zero pressure recovery for the flow downstream the flow separation point.
  - zero pressure fluctuation boundary condition at the vocal fold outlet;
- Constant flow rate at the vocal fold inlet
  - zero velocity fluctuation.
  
- Linearized 1D momentum and continuity equation around the mean state (mean vocal fold geometry)

$$\frac{\partial u}{\partial t} + U_0 \frac{\partial u}{\partial z} + u \frac{\partial U_0}{\partial z} + \frac{1}{\rho} \frac{\partial p}{\partial z} = 0$$

$$\frac{\partial h}{\partial t} + \frac{\partial (U_0 h)}{\partial z} + \frac{\partial (u H_0)}{\partial z} = 0$$



# Flow Pressure on the FSI interface

$$\frac{p(z,t)}{\rho} = \int_{z^{\text{sup}}}^{\bar{z}} \left( \frac{\partial}{\partial t} + U_0 \frac{\partial}{\partial z} + \frac{\partial U_0}{\partial z} \right) \left[ \frac{1}{H_0} \int_{z^{\text{inf}}}^{\bar{z}} \left( \frac{\partial}{\partial t} + U_0 \frac{\partial}{\partial z} + \frac{\partial U_0}{\partial z} \right) h dz \right] dz$$

$$= p_0(h) + p_1(h_t) + p_2(h_{tt})$$

$$p_0(h) = \int_{z^{\text{sup}}}^{\bar{z}} \left( U_0 \frac{\partial}{\partial z} + \frac{\partial U_0}{\partial z} \right) \left[ \frac{1}{H_0} \int_{z^{\text{inf}}}^{\bar{z}} \left( U_0 \frac{\partial}{\partial z} + \frac{\partial U_0}{\partial z} \right) h dz \right] dz$$

Flow-induced stiffness

$$p_1(h_t) = \int_{z^{\text{sup}}}^{\bar{z}} \left[ \frac{1}{H_0} \int_{z^{\text{inf}}}^{\bar{z}} \left( U_0 \frac{\partial}{\partial z} + \frac{\partial U_0}{\partial z} \right) h_t dz \right] dz + \int_{z^{\text{sup}}}^{\bar{z}} \left( U_0 \frac{\partial}{\partial z} + \frac{\partial U_0}{\partial z} \right) \left[ \frac{1}{H_0} \int_{z^{\text{inf}}}^{\bar{z}} h_t dz \right] dz$$

$$p_2(h_{tt}) = \int_{z^{\text{sup}}}^{\bar{z}} \frac{1}{H_0} \int_{z^{\text{inf}}}^{\bar{z}} h_{tt} dz dz$$

added mass,  
present even for  
zero flow rate

Flow-induced damping

Generalized Force

$$Q = Q_2 \ddot{q} + Q_1 \dot{q} + Q_0 q$$



# Eigenvalue problem

- Control equations

$$(M - Q_2)\ddot{q} + (C - Q_1)\dot{q} + (K - Q_0)q = 0$$

- Constant loss factor  $\sigma$

$$C = \sigma\omega M$$

- Eigenvalue Formulation

$$\begin{bmatrix} \dot{q} \\ \ddot{q} \end{bmatrix} = \begin{bmatrix} 0 & I \\ -(M - Q_2)^{-1}(K - Q_0) & -(M - Q_2)^{-1}(C - Q_1) \end{bmatrix} \begin{bmatrix} q \\ \dot{q} \end{bmatrix}$$

- Solve for eigenvalues/eigenvectors as a function of jet velocity/subglottal pressure
- Phonation onset occurs at the jet velocity for which the growth rate of one eigenvalue first becomes positive



# Solution Procedure

- Two-step procedure:
  - 1. Solve for steady state for a certain imposed subglottal pressure
  - 2. Solve the eigenvalue problem, checking for phonation onset. If no onset, repeat step 1 for a larger value of subglottal pressure. If onset, stop.
- For simplicity, the first step was skipped in this study.
  - Assume the mean steady-state vocal fold geometry does not change with subglottal pressure
  - Start with given steady-state (pre-phonatory) vocal fold geometry
  - Linear stability of the given pre-phonatory vocal fold geometry.



# Model Parameters Used

	Non-dimensional values	Physical value
Structural Damping Loss factor	0.1 (weak)	0.1
VF Thickness	1	7 mm
VF Depth	1.4	10 mm
VF Entrance Slop	45°	45°
Glottal Channel Gap	0.1428	1 mm
VF Density	1	1030 kg/m <sup>3</sup>
Young's Modulus	1	3 kPa
Scaling velocity	1	1 m/s
Flow Density	0.0012	1.2 kg/m <sup>3</sup>



# Effects of Flow-induced Stiffness

- Control equations

$$(M - \cancel{Q_2})\ddot{q} + (C - \cancel{Q_1})\dot{q} + (K - \cancel{Q_0})q = 0$$

- Effects of Flow-induced damping and inertia were studied and reported in another study\*.
- Effects of Flow-induced Stiffness:

$$M\ddot{q} + (K - Q_0)q = 0$$

\* Zhang, Z., Neubauer, J., Berry, D.A., (2007) "Physical mechanisms of phonation onset: a linear stability analysis of an aeroelastic continuum model of phonation," in review, J. Acoust. Soc. Am.



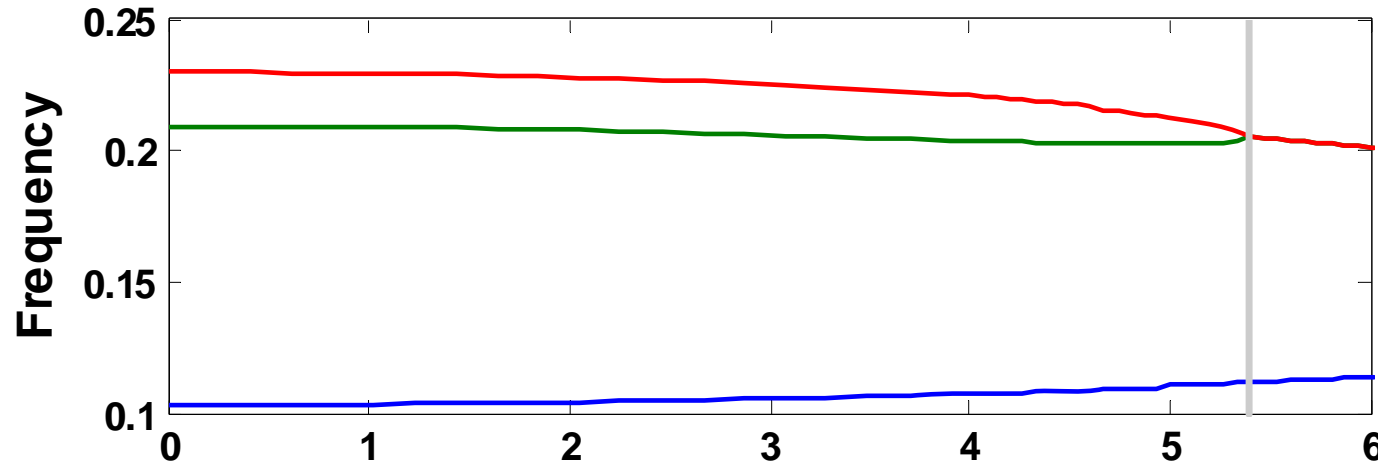
# Possible instabilities of $M\ddot{q} + (K - Q_0)q = 0$

- Static Divergence: zero-frequency instability
  - Flow-induced stiffness exceeds the structural stiffness
  - $Q_0$  is symmetric: the system is a conservative system, in which limit cycle solutions are impossible (see, e.g., Strogatz, 1994)
  - One-eigenmode instability
- Coupled-mode flutter: phonation onset
  - $Q_0$  is asymmetric, non-conservative system
  - Synchronization of two structural modes



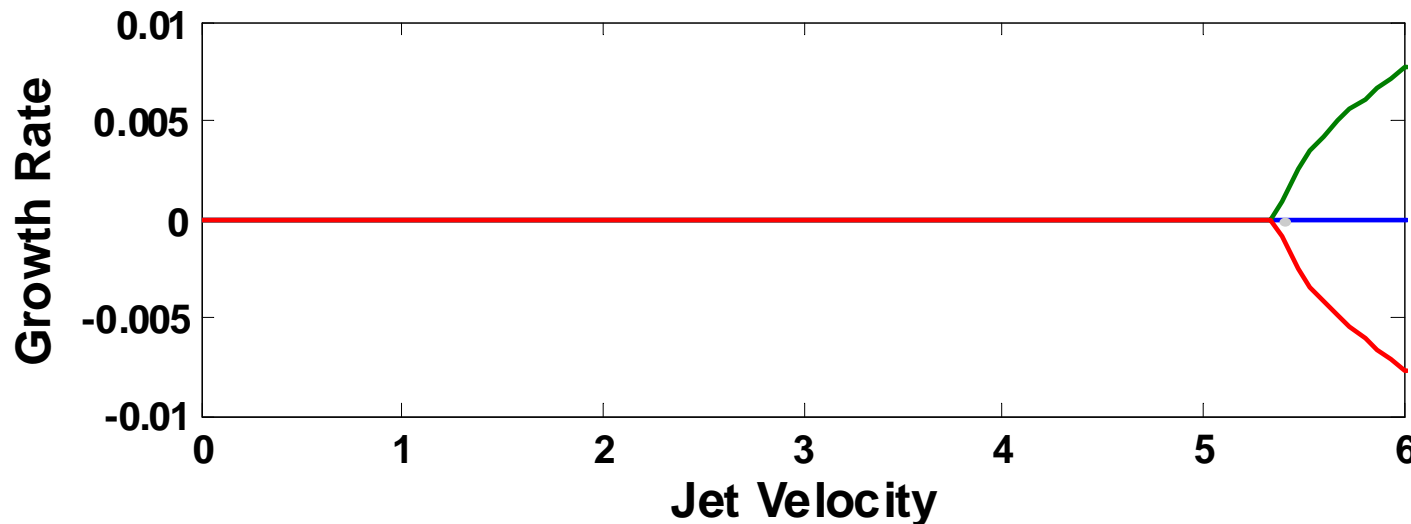
# Synchronization of two eigenmodes

$$M\ddot{q} + (K - Q_0)q = 0$$



Eigenmodes move toward each other due to the flow coupling effects

Two eigenmodes merge and lead to phonation onset



Similar behavior has been reported in the two-mass model (Ishizaka, 1981; 1986)



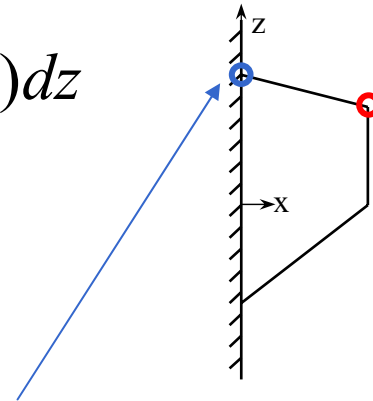
# Role of Flow separation in phonation onset

- The flow-induced stiffness matrix is asymmetric due to imposed flow separation

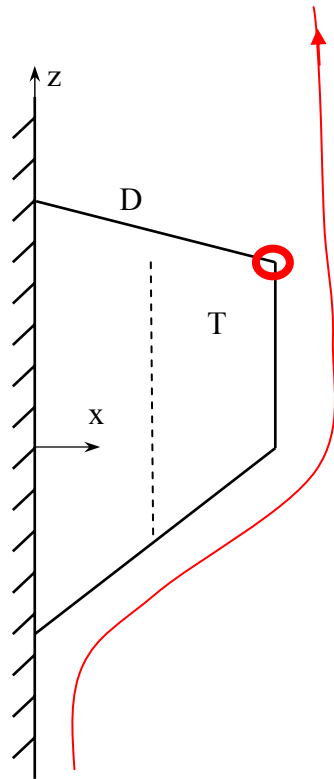
$$Q_0(i, k) = \frac{U_0^2}{H_0} \int_{z_{\text{inf}}}^{z_{\text{sup}}} \varphi_i \varphi_k dz - \frac{U_0^2(z_{sp})}{H_0(z_{sp})} \int_{z_{\text{inf}}}^{z_{\text{sup}}} \varphi_i(z_{sp}) \varphi_k(z) dz$$

$\varphi_i = x^m z^n$  is the basis function of the Ritz method

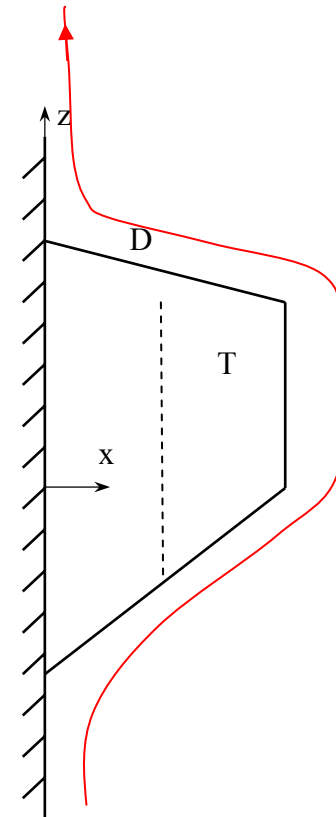
- $Q_0$  is symmetric only when  $\varphi_i(z_{sp}) = 0$  for all the basis functions.



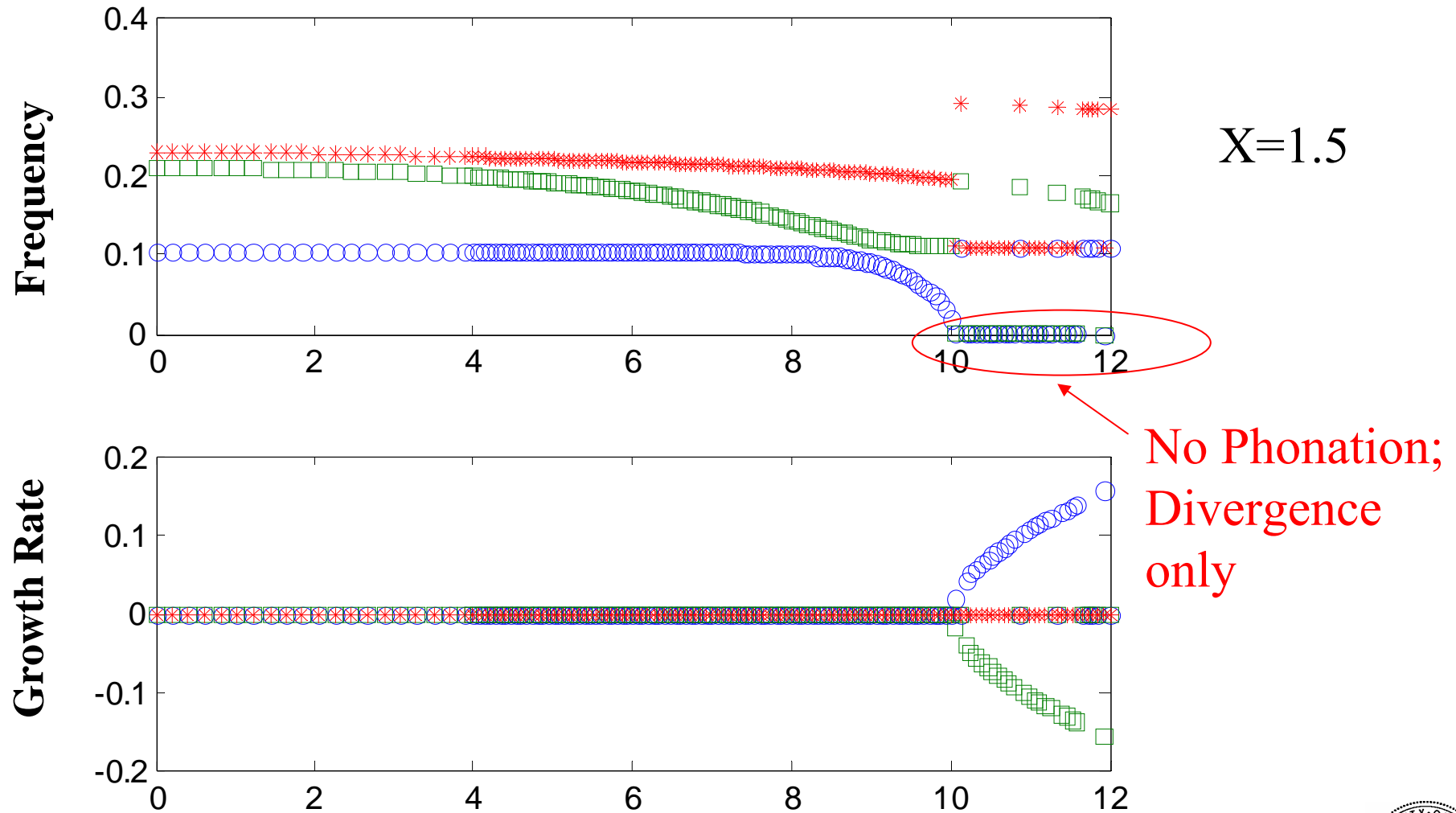
# Separated Flow



# Potential Flow



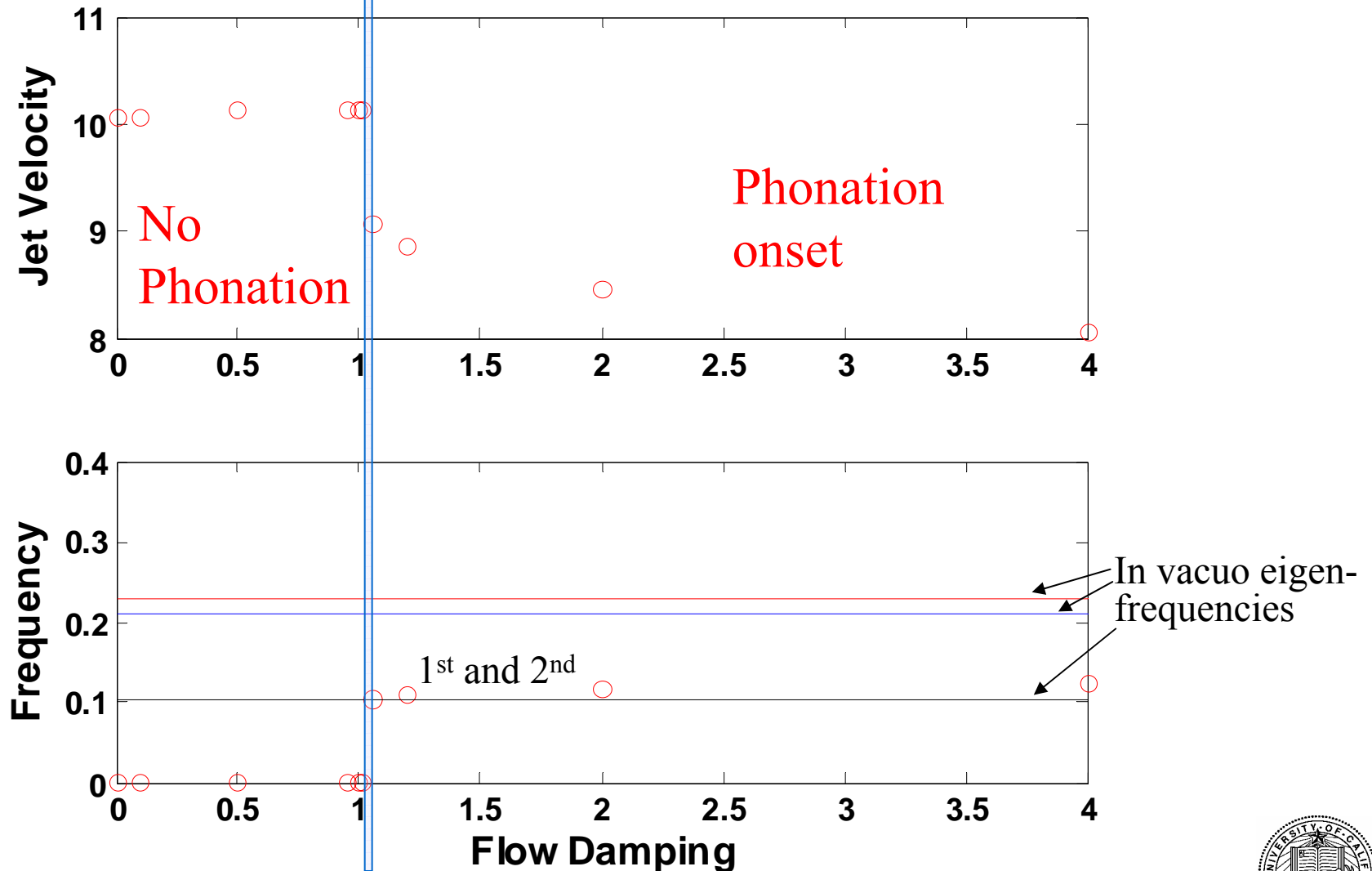
# Potential flow: No Flow Separation



- Any non-conservative mechanisms would contribute to the asymmetric nature of the  $Q_0$  matrix, and facilitate phonation onset:
  - e.g., sufficiently large viscous flow resistance



# Effects of Flow Damping without flow separation



# Effects of flow separation location

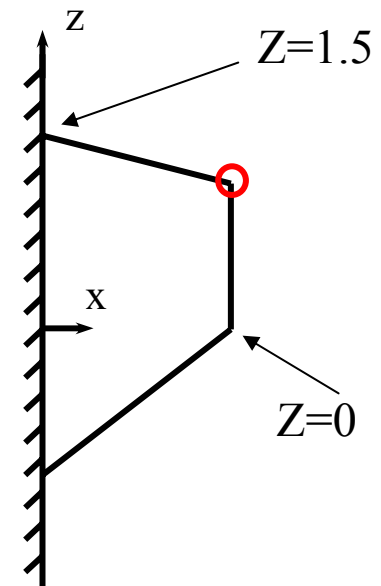
- Flow separation  $\rightarrow$  asymmetric flow-induced stiffness matrix  $Q_0 \rightarrow$  eigenmode-synchronization  $\rightarrow$  phonation onset
- Variation of the flow separation point may affect the degree of asymmetry of the  $Q_0$  matrix, and therefore affect the eigenmode-synchronization pattern.



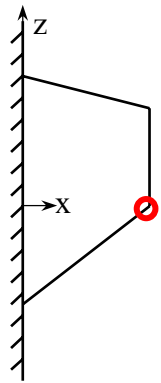
# Effects of flow separation location

Influence on:

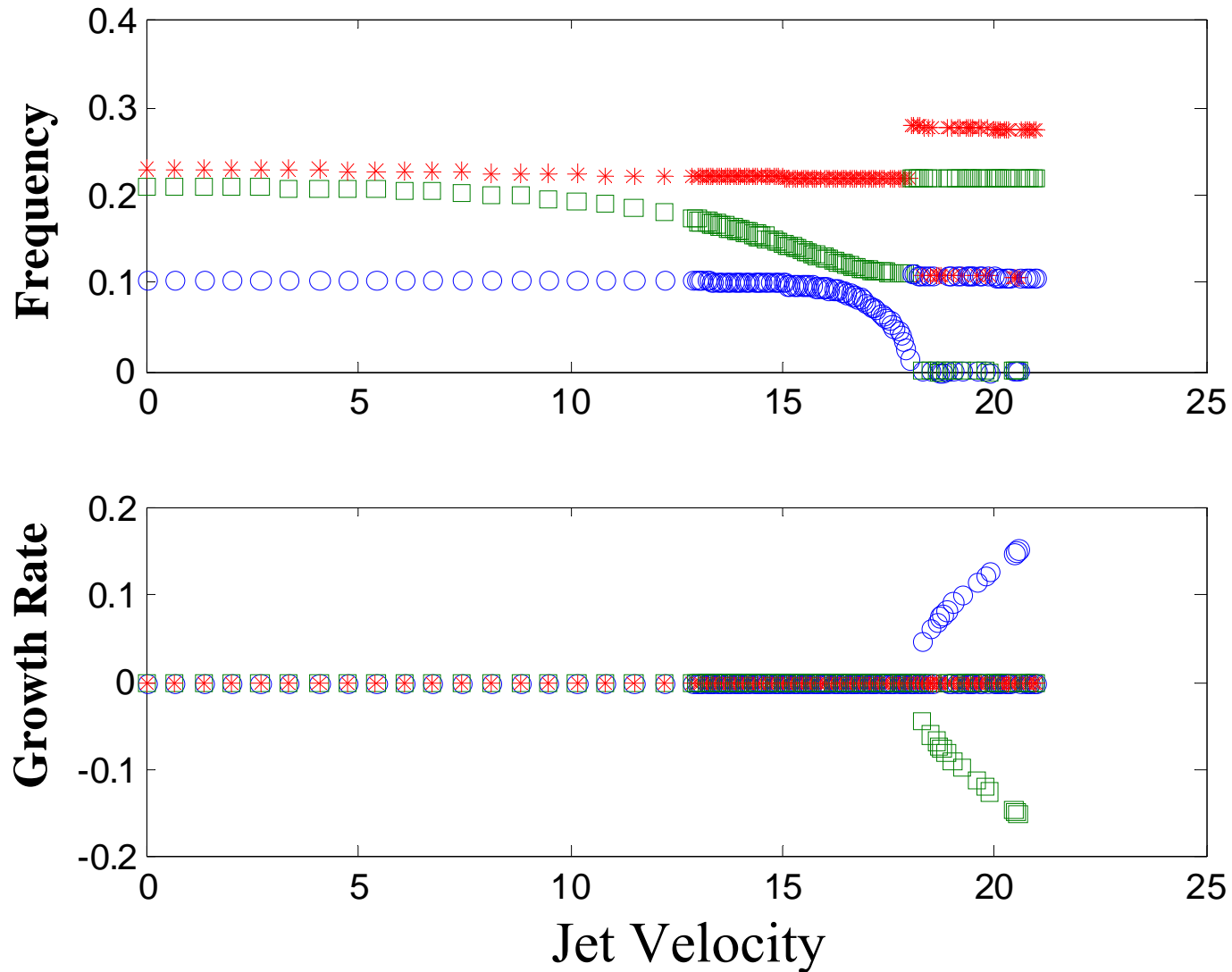
- Eigenmode synchronization pattern
- Phonation threshold
- Phonation frequency at onset
- Move the fixed flow separation point along the medial surface
  - $0 \leq Z_{sp} \leq 1.5$



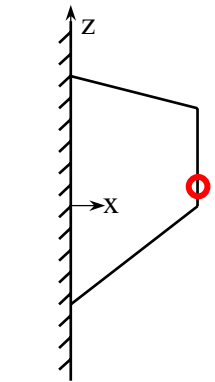
# Effects of Flow Separation Location



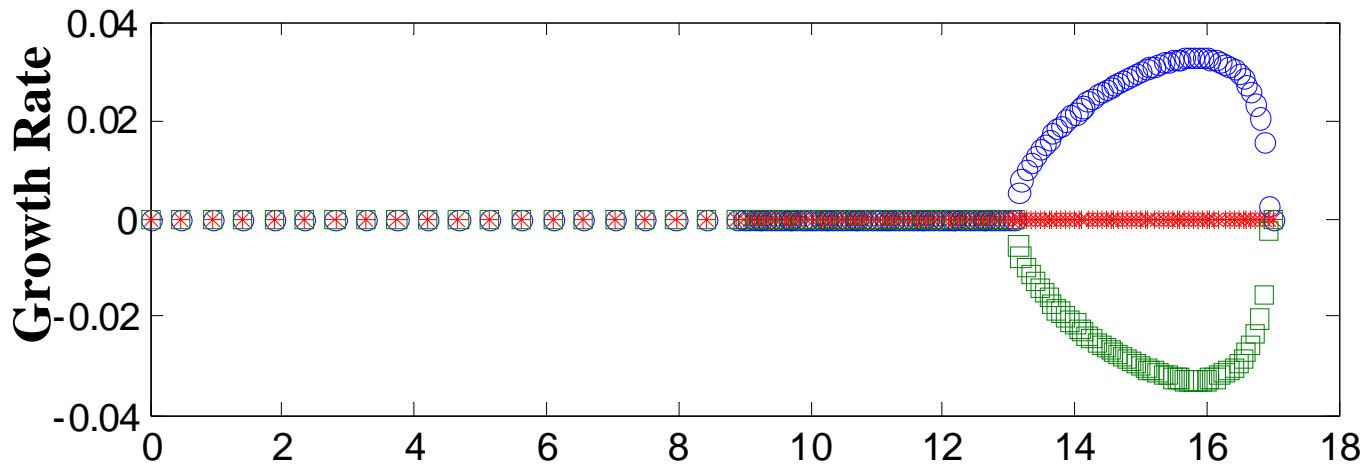
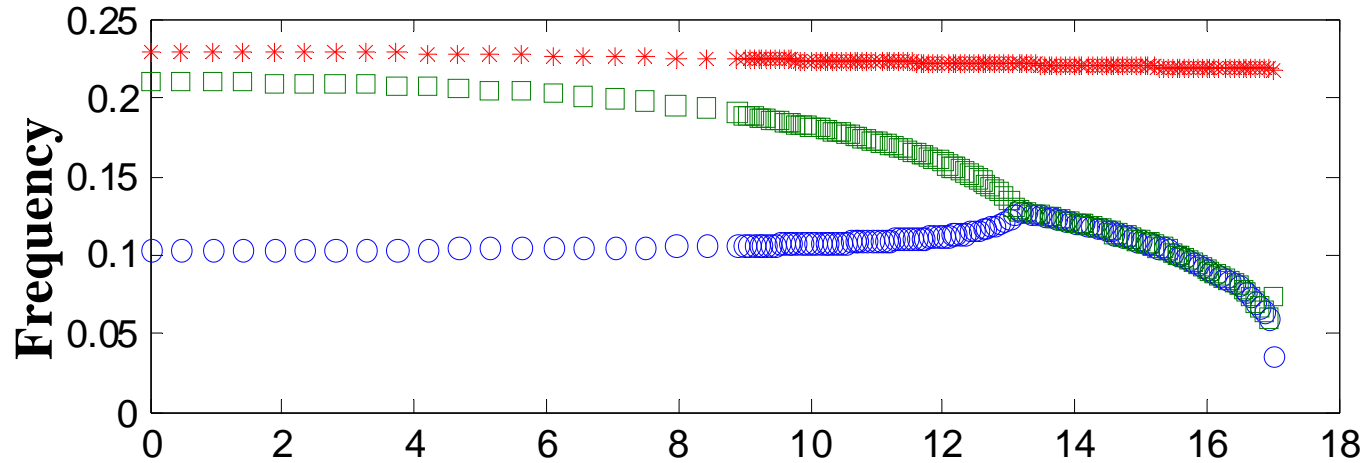
$$Z_{sp} = 0$$



# Effects of Flow Separation Location



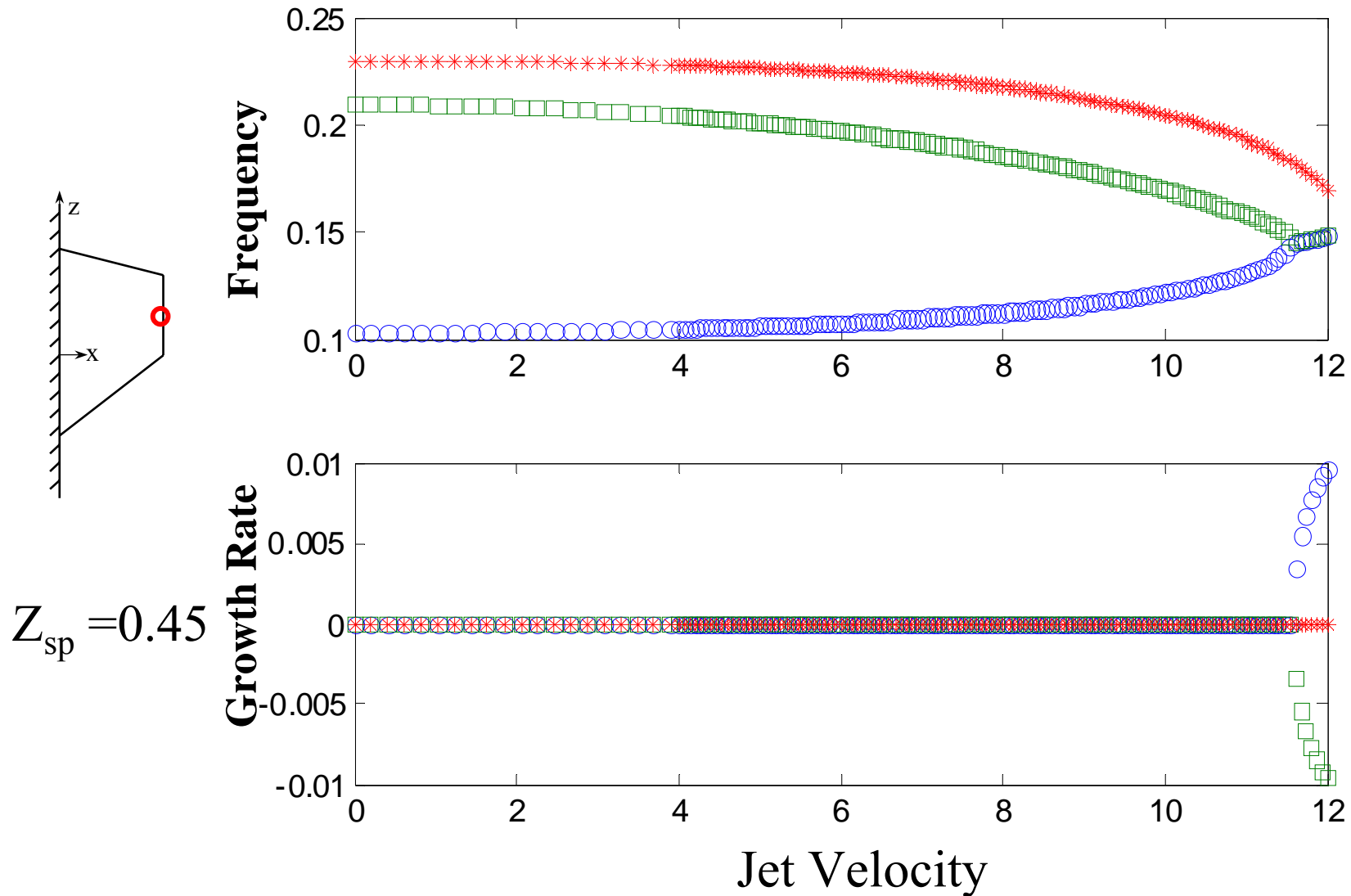
$$Z_{sp} = 0.1$$



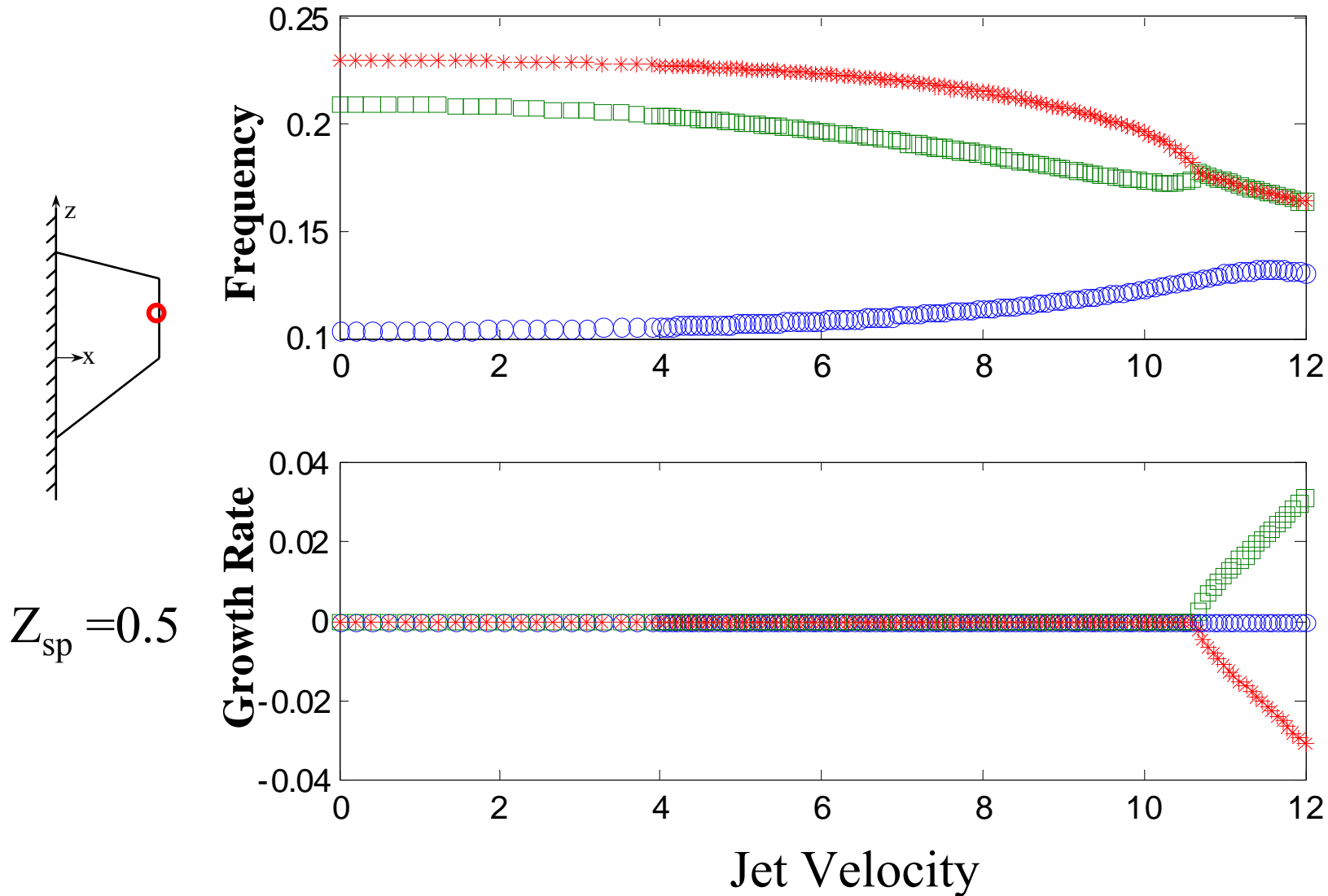
Jet Velocity



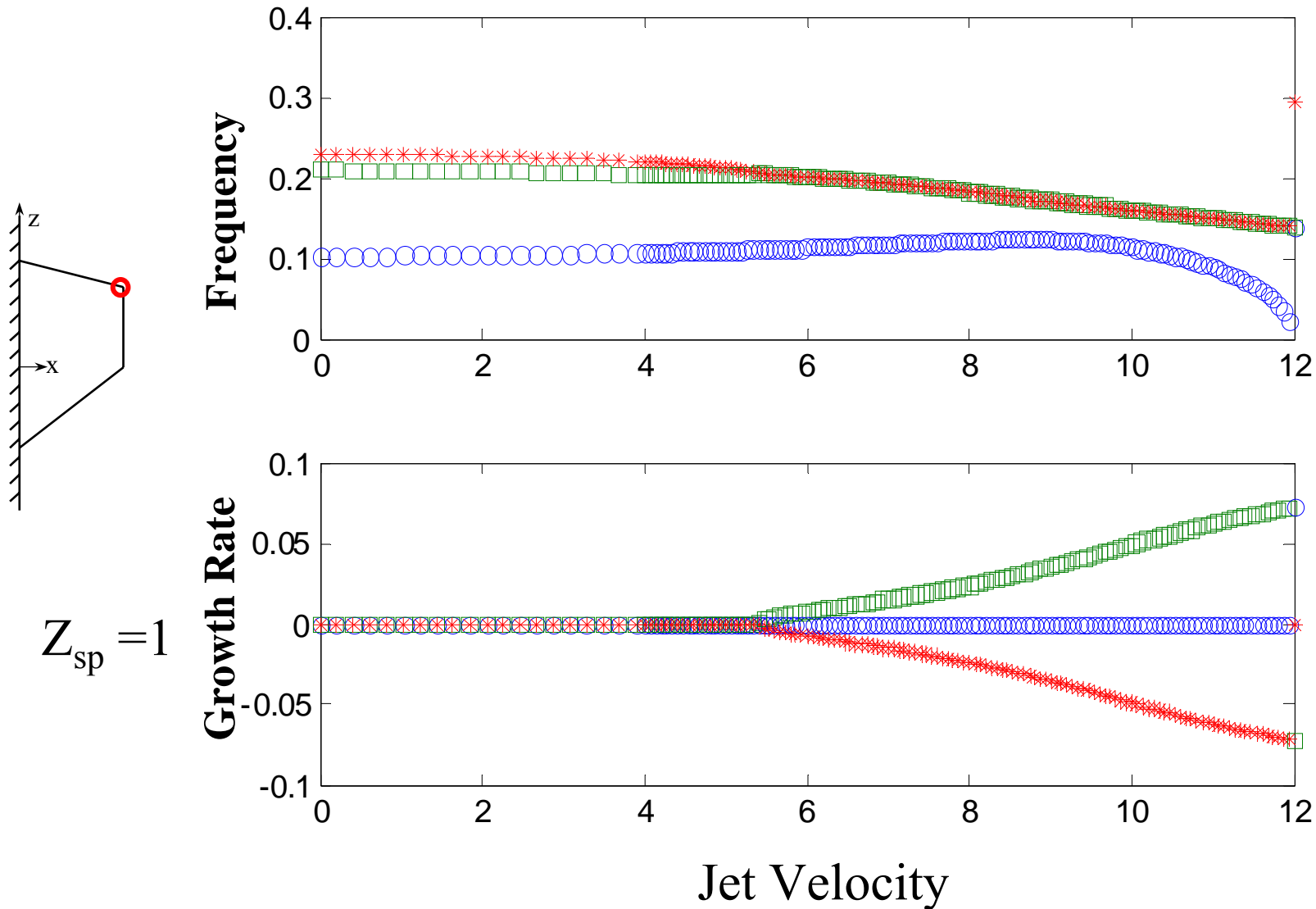
# Effects of Flow Separation Location



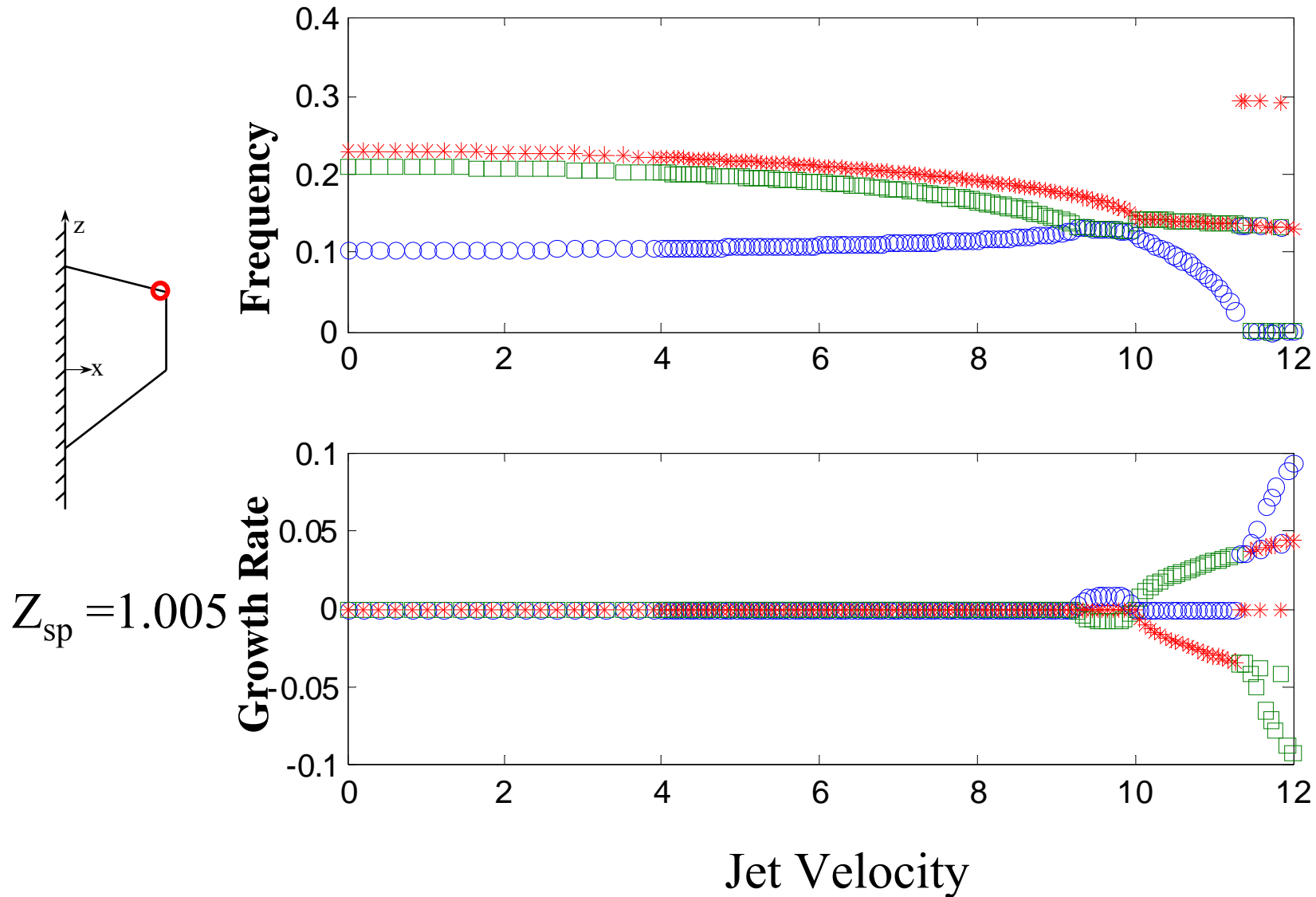
# Effects of Flow Separation Location



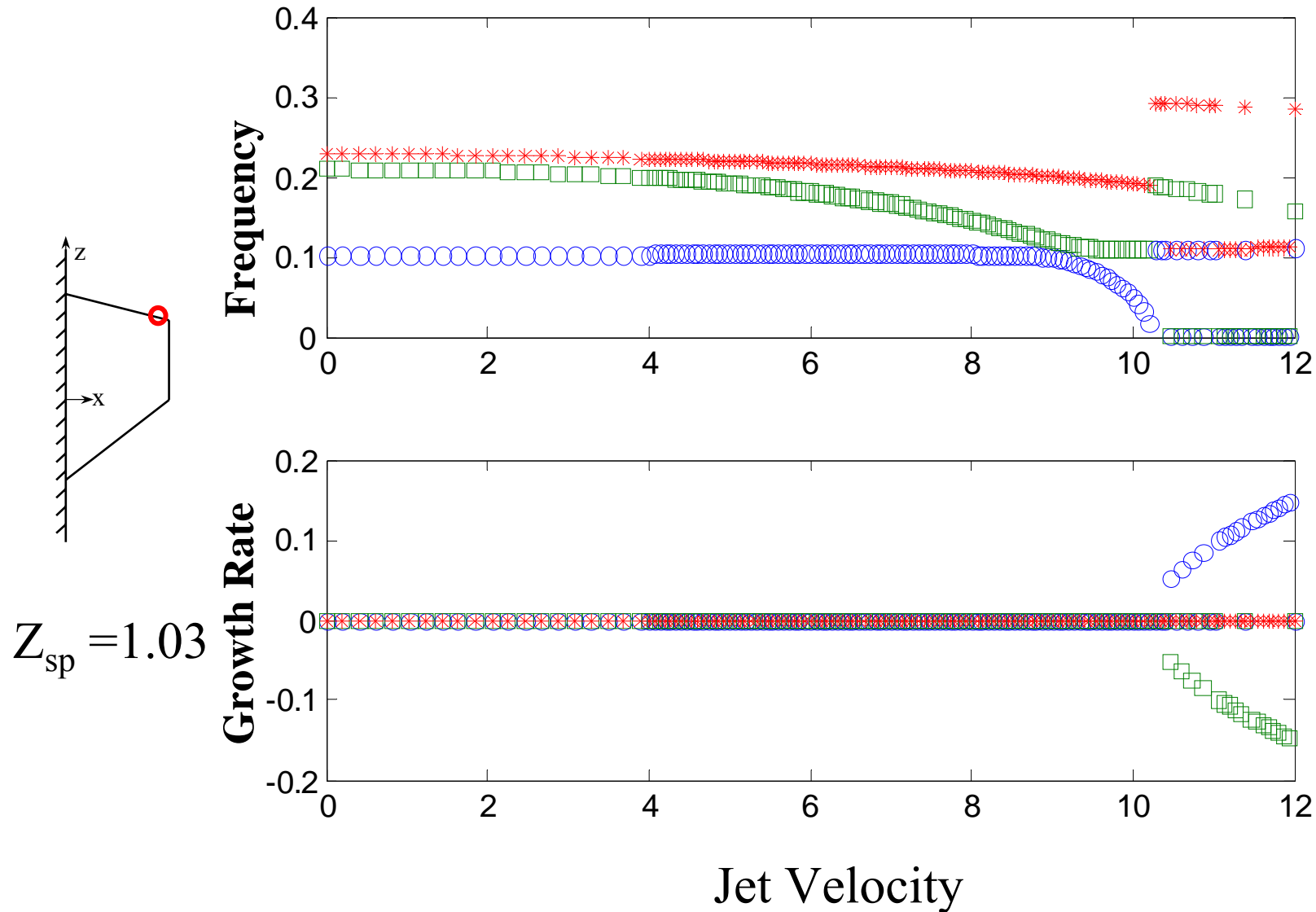
# Effects of Flow Separation Location



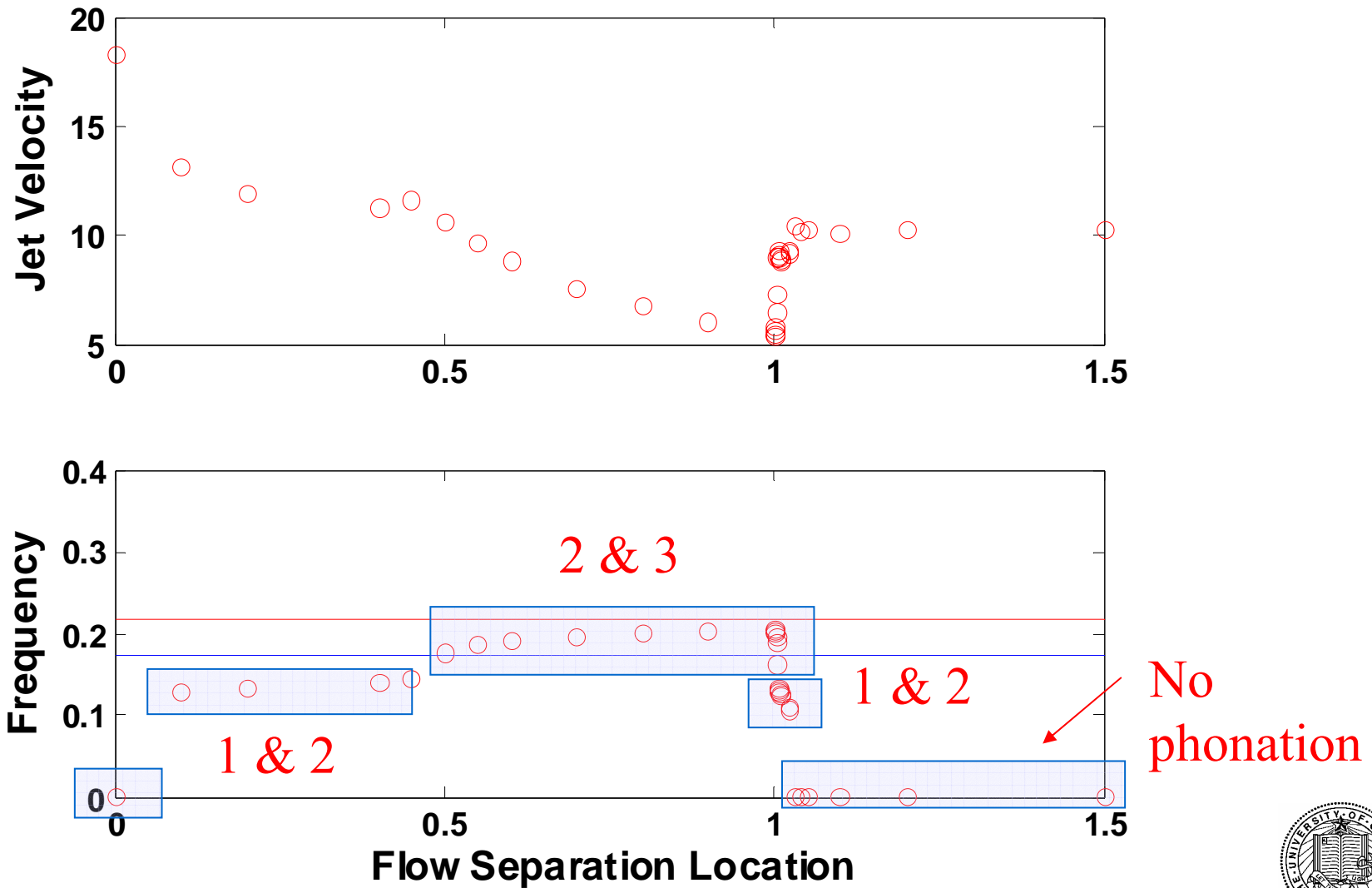
# Effects of Flow Separation Location



# Effects of Flow Separation Location



# Effects of Flow Separation Location



# Effects of Flow Separation Location - II

- Movement of the flow separation location changes the eigenmode interaction pattern (or movement of eigenmodes in the phase space), leading to different types of instabilities and eigenmode-synchronization patterns:
  - 1<sup>st</sup> & 2<sup>nd</sup>, 2<sup>nd</sup> & 3<sup>rd</sup>, and divergence.
- Eigenmode synchronization does NOT always occur between two closely-spaced in vacuo eigenmodes.



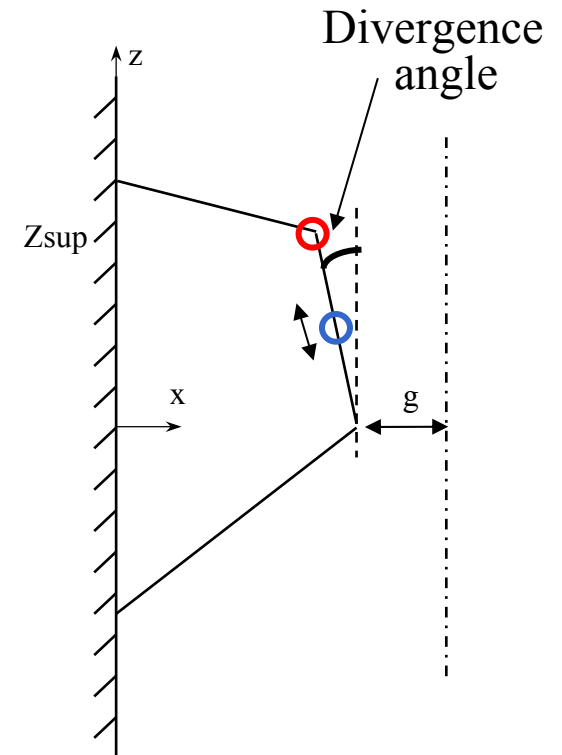
# Effects of glottal channel geometry

(convergent, straight, and divergent)



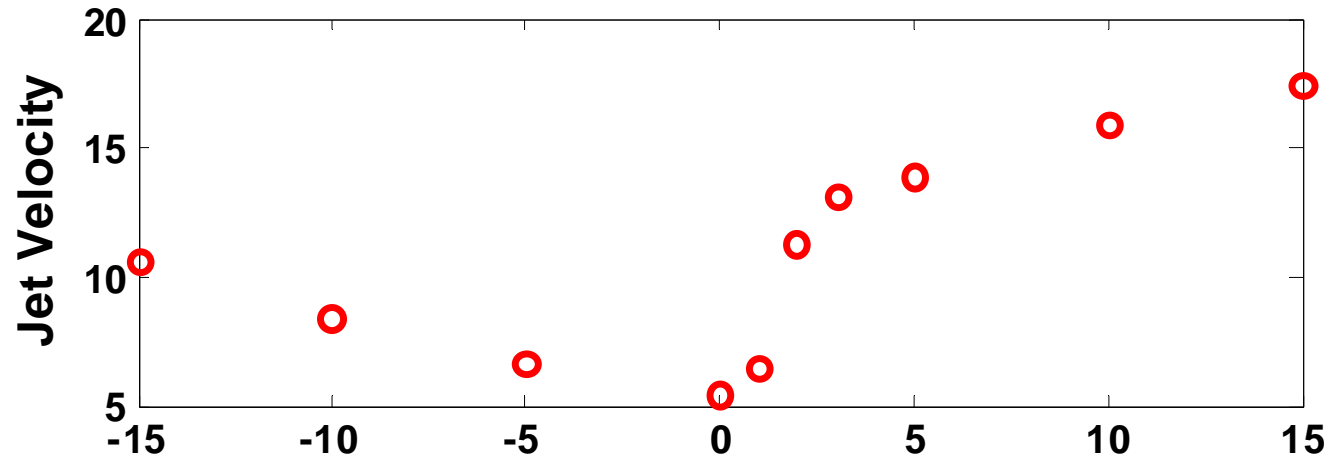
# Vocal fold with varying divergence angle

- Varying divergence angle
- Minimum glottal half-width ( $g$ ) remains constant
- Flow separation location determined in two ways:
  - 1. Fixed flow separation location at the superior edge of the medial surface.
  - 2. Flow separation location dependent of divergence angle.
    - Criterion for determination of the flow separation location (the  $A_{sp}/A_{min}=1.2$  criterion):
    - Convergent:  $Z_{sp}=Z_{min}$
    - Straight:  $Z_{sp}=Z_{sup}$
    - Divergent:
      - If  $Z_{sup} < Z_{min} * 1.2$ ,  $Z_{sp}=Z_{sup}$
      - Otherwise  $Z_{sp}=Z_{1.2}$ , where  $A(Z_{1.2})=1.2 * Z_{min}$



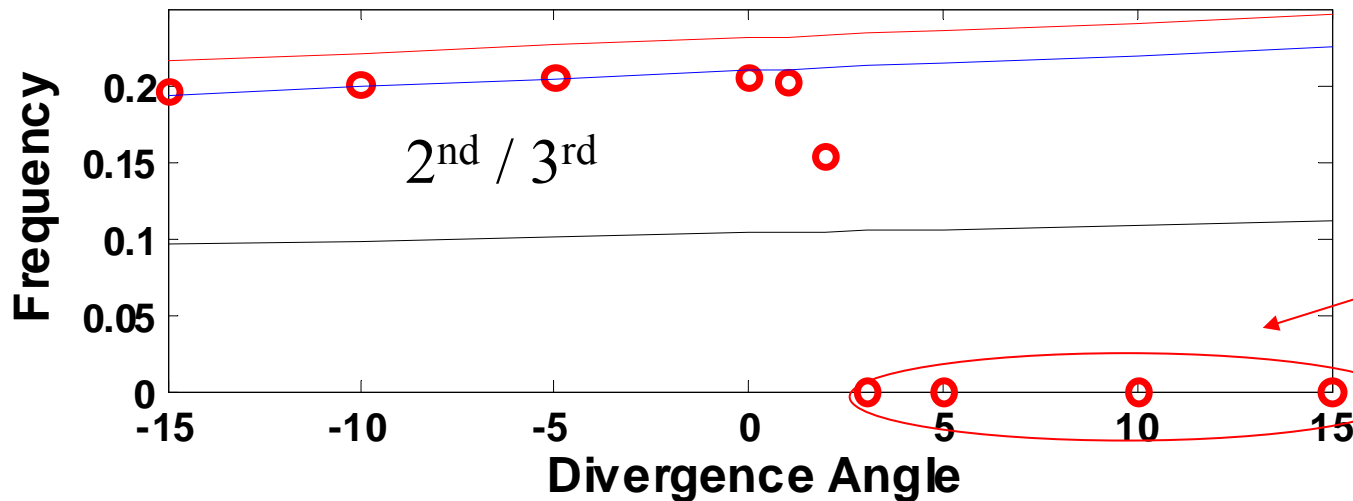
# Effects of Divergence Angle -- I

flow separation fixed at superior edge of medial surface



Minimum onset velocity occurred for straight glottal channel geometry.

Phonation frequency lower than the 2<sup>nd</sup> and 3<sup>rd</sup> in vacuo eigen-frequencies.

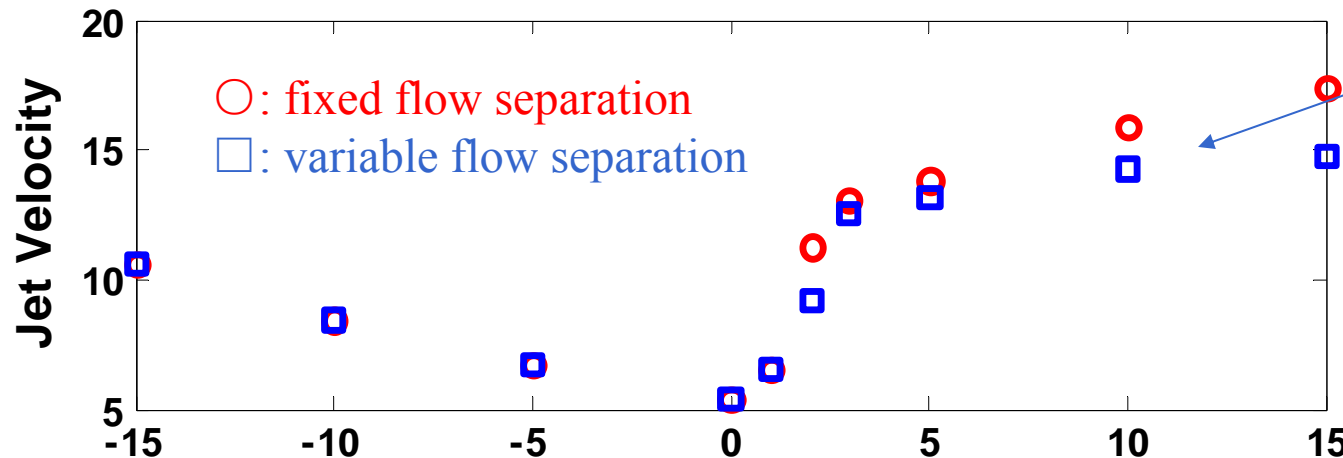


NO Phonation for divergent glottal channel and FIXED flow separation location



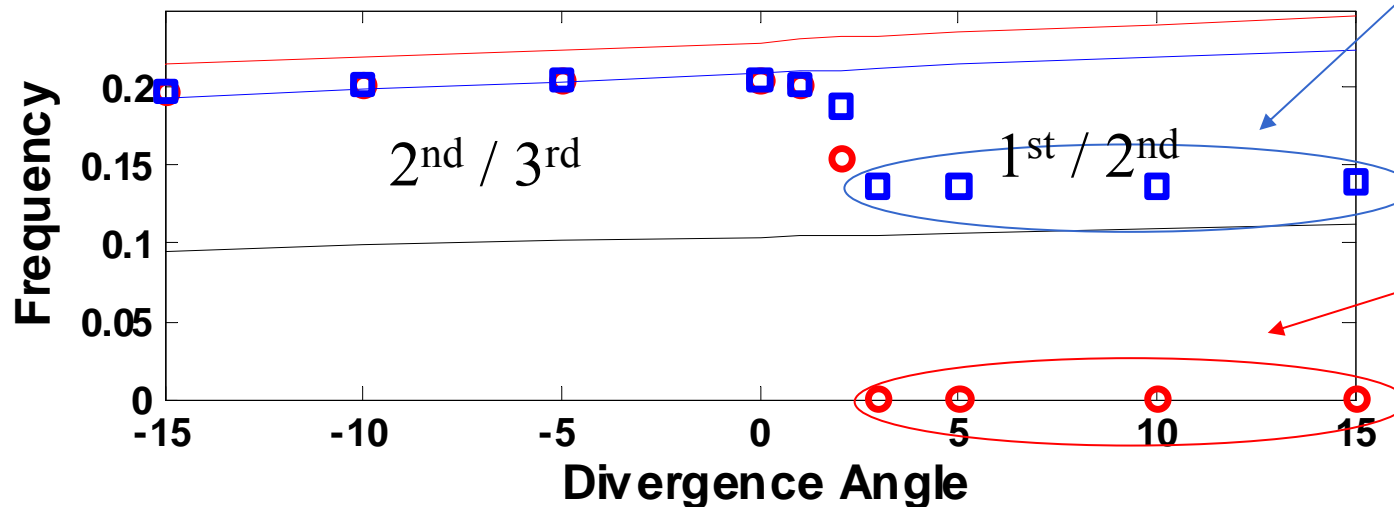
# Effects of Divergence Angle - II

## Flow separation location dependent on VF geometry



Reduced phonation threshold if flow separation location allowed to move

Phonation still possible for divergent glottal channel if flow separation location allowed to move



NO Phonation for divergent glottal channel and FIXED flow separation location



# Summary

- Phonation is due to the synchronization of two eigenmodes.
- Synchronization pattern affects the phonation threshold pressure, frequency, and vocal fold vibration
- Which eigenmodes synchronize is determined by the structure of the flow-induced stiffness matrix
  - many factors such as flow separation location, flow damping, and vocal fold geometry.
  - Synchronization does NOT always occur between two closely-spaced in vacuo eigenmodes
- Models needed to accurately predict the flow separation point
- At least two degrees-of-freedom need to be retained in reduced-order models of phonation
  - One-mass model or degenerate two-mass model can not capture the synchronization mechanism
  - Three degrees-of-freedom are recommended



# Future Work

- More realistic flow models:
  - Navier-stokes equations: effects of viscosity, presence of the boundary layer
- Subglottal and supraglottal tract acoustic loading
- Finite-amplitude analysis
- *Experimental verification.*

